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THE ROLE OF ANOMALOUS SATELLITE-FIXED ACCELERATIONS IN TOPEX/POSEIDON ORBIT MAINTENANCE*

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Shortly after the launch of TOPEX/Poseidon on 10 August 1992 orbit determination indicated orbital decay levels ~60 times larger than could be explained by atmospheric drag. Outgassing, a complex process of molecular releases from satellite non-metallic parts, was the most likely source of these early decay rates. The high decay levels steadily declined during the first six weeks while a planned sequence of six orbit adjust maneuvers placed the satellite in the operational orbit to precisely overfly a predetermined repeat ground track. At the same time, on-going orbit trend analysis revealed the presence of residual along-track forces comparable to atmospheric drag which clearly exhibited a body-fixed origin. These anomalous forces cause either orbital decay or boost, depending on the satellite attitude and solar array articulation mode. As such, these along-track forces can either add to, or oppose, the resident orbital decay due to drag.

Orbit maintenance maneuver design was expected to depend primarily on effective predictions of atmospheric drag, but now also depends equally on reliable predictions of the anomalous along-track forces. The basic behavior of the anomalous forces has been established, evolving from quite unknown character early in the mission to reasonably predictable by mid-August 1993. This paper describes the method used to estimate the anomalous forces and presents an empirical prediction model for each of the four satellite yaw control modes, also reflecting the significant influences of a newly-adopted operational strategy to bias solar array pointing for battery lifetime enhancement.

INTRODUCTION

TOPEX/Poseidon was successfully launched by an Ariane 42P from French Guiana on 10 August 1992. The primary goals of this joint US/French mission are to study ocean circulation and its interaction with the atmosphere, to better understand climate change; to improve knowledge of heat transport in the ocean; to model ocean tides; and to study the marine gravity field. To accomplish these objectives requires determination of ocean surface height to an accuracy of 13 cm utilizing a combination of satellite altimetry and precision orbit determination¹ based primarily on laser ranging measurements. These objectives are to be accomplished over a primary mission lifetime of three years, with a possible two-year extension.

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The Jet Propulsion Laboratory (JPL) is responsible for TOPEX/Poseidon mission operations, including operational navigation. Major navigation functions include all maneuver design, evaluation, and related orbit analysis. Operational orbit determination support is provided to JPL by the Flight Dynamics Facility at the Goddard Space Flight Center (GSFC/FDF) using tracking data acquired via the NASA Tracking and Data Relay Satellite System (TDRSS). Detailed interfaces and procedures for exchanging maneuver and orbit determination data between JPL and the GSFC/FDF were established and thoroughly tested prior to launch to demonstrate that all performance requirements could be satisfied.²

ORBIT MAINTENANCE REQUIREMENTS

A planned sequence of six orbit adjust maneuvers³ began soon after launch to precisely place the satellite in a near-circular frozen orbit at an inclination of $\sim 66^\circ$ and an equatorial altitude of ~ 1336 km. During the maneuvering process, the orbit ground track was precisely aligned with a reference ground track which repeats every 127 revolutions over ~ 9.9 days, while also overflying single US and French altimeter verification sites. This sequence was completed on 21 September 1992, 42 days after launch.⁴

Periodic maneuvers maintain the operational orbit and attendant ground track. The specified control and maneuver scheduling constraints require that:

- 1) 95% of all equatorial crossings are contained within a 2-km longitude band at each orbit node,
- 2) 95% of all altimeter verification site overflights are within ± 1 km during the initial verification phase (at least first six months),
- 3) maneuver spacings are at least 30 days, time-phased to occur near the boundary of pre-determined 9.9-day ground track repeat cycles to limit interference with precision orbit determination, and
- 4) the burn occurs over land to preclude interruption of ocean altimetry.

Maneuver design also reflects satellite health and safety concerns by satisfying telecommunications, thermal, and power constraints when selecting the maneuver location near the cycle boundaries.

SATELLITE YAW CONTROL MODES AND SOLAR ARRAY POINTING STRATEGY

YAW CONTROL MODES

TOPEX/Poseidon is a three-axis stabilized satellite (Figure 1) with the altimeter boresight always pointed normal to the reference ellipsoid. At the same time, near-continuous sinusoidal yaw steering about the local nadir and solar array (SA) pitching combine to maintain the SA pointed near the sun for power optimization. The yaw steering strategy is used continuously except when $-15^\circ < \beta' < 15^\circ$, where β' is the angle between the orbit plane and the sunline (e.g., Figures 2c and 11). When β' is near these angular limits a fixed yaw attitude is utilized to avoid excessive yaw rates. The satellite is positioned at a zero yaw angle ($\psi_s = 0^\circ$) when $0^\circ < \beta' < 15^\circ$ (flying forward), whereas $\psi_s = 180^\circ$ is utilized when $-15^\circ < \beta' < 0^\circ$ (flying backward). Accordingly, a yaw flip maneuver is required near $\beta' = 0^\circ$ to keep the SA on the sunlight side of the satellite. The β' angle passes through zero once every ~ 56 days as the satellite orbit node regresses $\sim 2.2^\circ/\text{day}$ and the earth moves in its orbit $\sim 1^\circ/\text{day}$.

The sinusoidal yaw steering algorithm is described in Eq. (1), where $\Omega = 90^\circ + \nu - \nu_{Sun}$, ν_{Sun} is the angle of the projection of the earth-sun line into the orbit plane measured from perigee, and ν is the true anomaly of the satellite. Solar noon is defined by $\Omega = 90^\circ$.

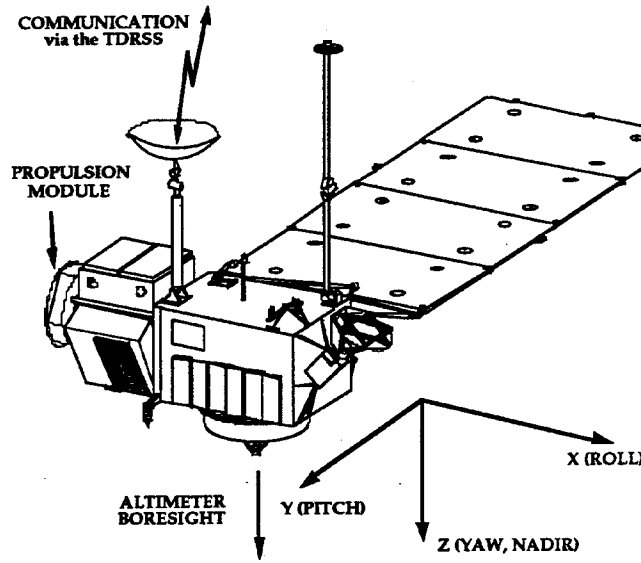


Figure 1 TOPEX/Poseidon Satellite

$$\begin{aligned}
 \beta' \geq 15^\circ: \quad \psi_s &= 90^\circ + (90^\circ - \beta') \cos \Omega && \text{(positive yaw steering)} \\
 0 < \beta' < 15^\circ: \quad \psi_s &= 0^\circ && \text{(flying forward)} \\
 -15^\circ < \beta' < 0: \quad \psi_s &= 180^\circ && \text{(flying backward)} \\
 \beta' \leq -15^\circ: \quad \psi_s &= -90^\circ - (90^\circ + \beta') \cos \Omega && \text{(negative yaw steering)}
 \end{aligned} \tag{1}$$

SOLAR ARRAY POINTING STRATEGY

In each yaw mode, the original plan⁵ was to continuously pitch the solar array so that it was as nearly normal to the sun as possible. This strategy required that the solar array pitch angle, γ , be set to γ_{opt} , where

$$\gamma_{opt} = 180^\circ + \tan^{-1} \left[\frac{\sin \Omega \cos \beta'}{\cos \psi_s \cos \Omega \cos \beta' - \sin \psi_s \sin \beta'} \right] \tag{2}$$

When the SA cell side normal parallels the satellite +X direction, $\gamma \equiv 0$. Provision was made to bias the solar array from its optimal direction by an angle B . In each yaw mode the solar array pitch angle is controlled to $\gamma_{opt} + B$, where $B > 0^\circ$ is referred to as a *leading bias*. The resulting pitch angle as a function of orbit angle Ω and yaw control mode is given by Eq. (3).

$$\begin{aligned}
 \gamma &= 180^\circ + B - \Omega && \text{for } \psi_s = 180^\circ \text{ (fixed yaw, fly backward)} \\
 \gamma &= \Omega + B && \text{for } \psi_s = 0^\circ \text{ (fixed yaw, fly forward)} \\
 \gamma &\approx 180^\circ + B - (90^\circ + \beta') \sin \Omega && \text{(negative yaw steering, } \beta' < -15^\circ) \\
 \gamma &\approx 180^\circ + B - (90^\circ - \beta') \sin \Omega && \text{(positive yaw steering, } \beta' > 15^\circ)
 \end{aligned} \tag{3}$$

The equations for yaw steering are approximate and are shown only to characterize the *windshield wiper* motion of the array in those control modes.

Shortly before launch, a plan was adopted to use the SA pitch bias B to limit the peak battery charging currents to 20 Amperes. A $+55^\circ$ value was initially employed on 28 August 1992; it was refined to 57.5° on 12 September 1992. More recently, on 27 July 1993, the bias was reduced to 53° to compensate for declining solar array output.

While this strategy seems to be lengthening the battery lifetime, it introduces forces normal to the sunline that are not included in the radiation pressure models adopted for joint use by JPL and the GSFC/FDF. These unmodeled forces introduce errors affecting the ground track that depend on the yaw control mode. Understanding and effectively modeling these effects are an important part of the orbit maintenance process and a major focus of this paper.

INITIAL MEASURES OF AN ANOMALOUS FORCE

Pre-launch studies⁶ established atmospheric drag as the major non-gravitational perturbation affecting the satellite ground track, even though the orbital altitude is relatively high at ~ 1336 km. This sensitivity to drag is a consequence of the stringent ± 1 -km ground track control requirement. Atmospheric drag causes decay in the orbit semi-major axis, resulting in an eastward drift of the satellite ground track that will eventually travel outside the established control boundary. Periodic maneuvers maintain the ground track inside the control boundaries by removing the accumulated orbital decay with an increase in semi-major axis.* The frequency of maneuvers depends on both the drag level and on the accuracy of drag predictions. Effective drag predictions require long-term predictions of atmospheric density derived from predictions of solar and geomagnetic activity.⁷ The Jacchia-Roberts atmospheric density model^{8,9} effectively accounts for daily variations in solar and geomagnetic activity, although none of the currently available density models reflect flight data at the TOPEX/Poseidon altitude. A faithful representation of the satellite variable mean area (VMA)⁶ is also needed. The VMA defines the orbital-average area as a function of β' based on a precise projected area model developed by Fairchild Space.¹⁰

To develop an initial trend in the observed orbital decay behavior, the GSFC/FDF first estimated daily drag acceleration multipliers using tracking data acquired via the TDRSS. The daily drag multiplier is $(1 + \rho_1)$, where $\rho_1 = 0$ indicates nominal drag. All unmodeled along-track accelerations are then arbitrarily absorbed as drag without necessarily declaring that these effects are actually due to drag. Figure 2(a) shows the trend in daily $(1 + \rho_1)$ estimates from launch through the end of 1992, indicating an exponential decline from an initial post-launch value of ~ 60 times nominal drag to near-nominal levels by late September 1992 when operational orbit conditions were finally achieved. Figure 2(b) shows the corresponding rate of change in semi-major axis (Eq. 4)^{11,12} varied between -200 and $+25$ cm/day, equivalent to along-track acceleration levels of ~ 1.2 to 8.3 nanometers/sec² (or forces between ~ 3 and 20 μ N).

$$\left(\frac{da}{dt}\right)_{\text{drag}} = -\frac{(1 + \rho_1)\rho AC_D \sqrt{\mu a}}{m} \left[1 - \frac{\omega_e \cos i}{n}\right]^2 \quad (4)$$

where $(1 + \rho_1)$ is the drag scale factor ($\rho_1 = 0$ indicates nominal drag), ρ is the average daily atmospheric density, A is the average per orbit reference drag area variation with β' (VMA), C_D is the satellite drag coefficient, μ is the earth gravitational constant, a is the orbit mean semi-major axis, ω_e is the earth rotation rate, i is the orbit inclination, and n is the orbit mean motion.

Normal outgassing is the most likely explanation for the initially high orbital decay rates. The accelerations associated with outgassing arise from imbalances between gas losses through different sides of the satellite. A differential loss of 0.4 grams, at 500 m/s velocity, will produce a force of 2.4 μ N, resulting in a 1 nm/s² acceleration. If aligned against the velocity vector, this force would cause ~ 18.5 cm/day decay in the mean semi-major axis. An equal force could be generated by only five puffs *per day* from a typical aerosol inhaler.

* Typical maneuver ΔV magnitudes are ≈ 5 mm/s, raising the semi-major axis by ≈ 10 meters.

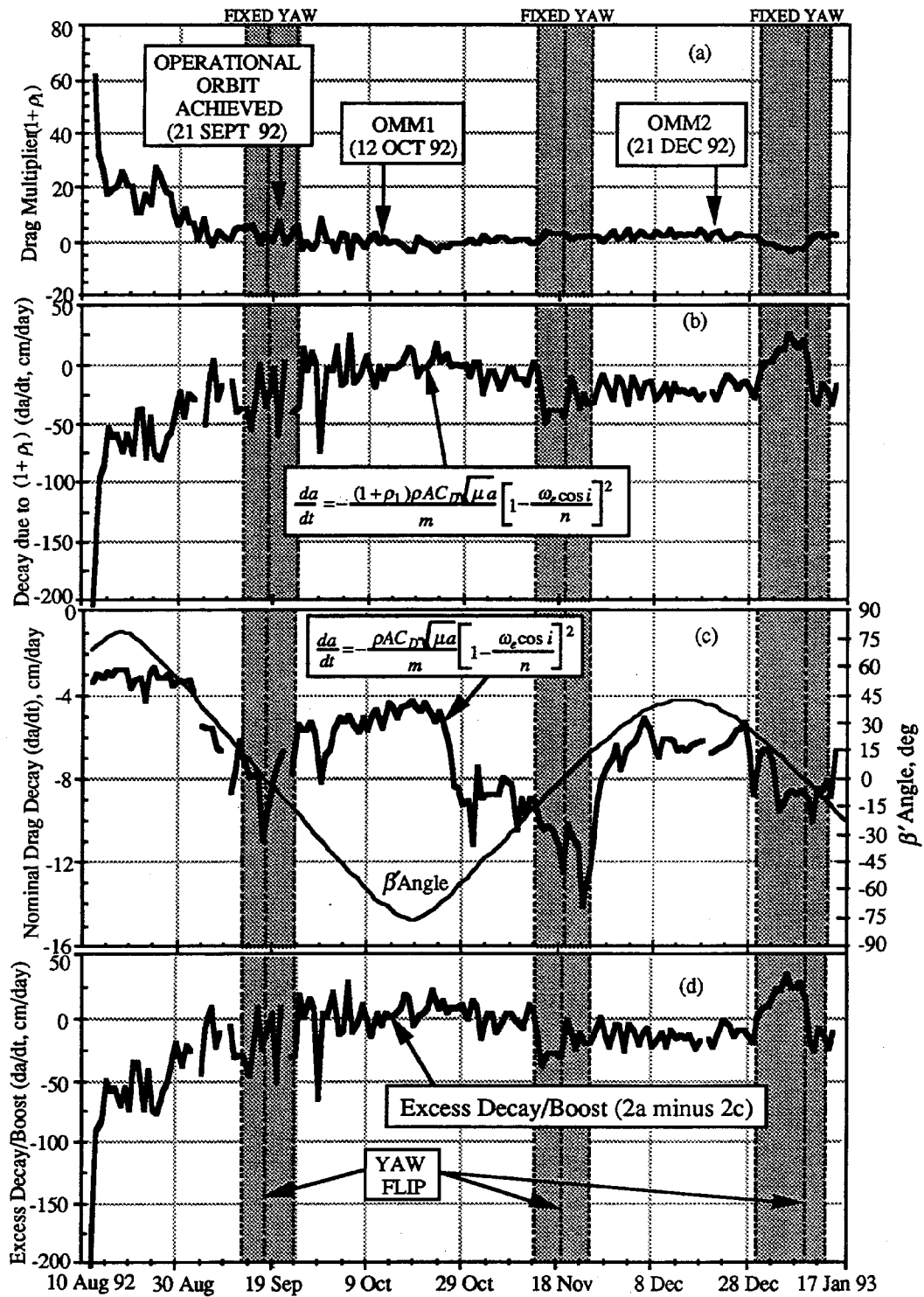


Figure 2(a) Atmospheric Drag Multiplier, (b) Equivalent Total Orbital Decay/Boost, (c) Orbital Decay Due to Observed Solar Activity, and (d) Excess Decay/Boost Rates.

The orbital decay rates due to drag, estimated by Eq. (4) with $\rho_1 = 0$, vary between ~ 3 and ~ 14 cm/day (Figure 2c). These decay rates are much lower than those implied by the $(1 + \rho_1)$ estimates, even after September 1992 when most of the suspected outgassing had subsided. The distinct long-term systematic signature is due to the VMA variation with β' , whereas the more frequent variations are caused by changes in the average daily atmospheric density with observed solar activity.

Estimates of orbital decay (or boost) rates in excess of nominal drag (Figure 2d) were obtained by removing the drag decay rates in Figure 2(c) from the total decay/boost rates in Figure 2(b). Figure 2(d) clearly shows excess *boost* rates during negative yaw steering, while excess *decay* rates prevail during periods of positive yaw steering. Also apparent are abrupt changes in the excess boost/decay levels upon entering and exiting periods of fixed yaw and after a yaw flip. These variations are consistent with a combination of satellite body-fixed forces and the forces arising from offsetting the SA pitch angle by B .

While this evaluation technique helped establish the presence of anomalous forces, it incorrectly assumes the forces are always proportional to the satellite drag acceleration. Instead, a separate estimate of an along-track force acting in the presence of nominal drag is needed to isolate the anomalous force. Before presenting these results, we will first identify potential sources of the anomalous force and provide some physical interpretations.

SOURCES OF OBSERVED ANOMALOUS FORCES

INTRODUCTION

There have been extensive efforts to model the forces acting on the satellite. The Precision Orbit Determination Team (PODT) at GSFC has developed a *box and wing* model of the satellite^{13,14} and later tuned it using laser ranging data and dual-frequency Doppler measurements from the French DORIS tracking system.¹⁵ This modeling effort will be continually refined.

In some cases, the force models in the Navigation Team (NAVT) software are considerably simpler than their PODT counterparts in order to minimize the computational burden while providing needed accuracy for predicting the long-term ground track behavior. Drag and sunline forces are based on the VMA,⁶ while the gravity field is a 20x20 truncation of the GEMT3 model.* The PODT requires a much larger degree and order model (up to 70x70). The NAVT models were entirely adequate to the task *until* the decision to offset the SA by angle B was made shortly before launch.

Several potential sources of the *anomalous* forces were identified:

- 1) Outgassing from the interior of the satellite. These body-fixed forces are expected to be constant in direction and to decline with time as volatiles *bake out*. Thus all body-fixed components should decay simultaneously. Outgassing is believed to be the primary source of the early *large* orbital decay forces.
- 2) Curling of the solar array due to temperature differences between front and back produces a force in the satellite +Y direction. This force is primarily dependent on reflected sunlight and is proportional to the curl angle.
- 3) There is a net thermal force due to radiation from the satellite surfaces. Such forces are treated as body-fixed, but can vary with sun angle and yaw mode.
- 4) Unmodeled (in the NAVT software) forces arising from reflections from the offset solar array. The PODT carefully models these forces.

* The NAVT used a 20x20 truncation of GEMT3 from launch until 27 July 1993, when it was replaced by a 20x20 truncation of JGM2 at the same time the SA pitch offset angle B was reduced from 57.5 to 53 deg.

Other sources considered but not discussed here include outgassing from the satellite blankets (after heating) and the (unmodeled) high-gain antenna used to communicate through the TDRSS. The possibility of leaky thrusters was quickly rejected in favor of an outgassing model for the early large declining force.

Our brief analysis is presented in two parts. The first part demonstrates the effect of body-fixed forces in the different yaw modes. The second part estimates the effect of the unmodeled reflections from the solar array in the different yaw modes. By such analytic modeling the effect of changing sun angle, solar array offset, and other factors can be predicted and serve to verify and help interpret inflight observations.

EFFECTS OF SATELLITE BODY-FIXED FORCES

Both pre-launch studies⁶ and perturbation analyses have shown that only forces along the velocity vector are important in the long-term behavior of the ground track. Secular perturbations in eccentricity and argument of periape can arise from once-per-orbit forces, both radial and along track, but these effects need not be considered here.

Shapiro¹⁶ was the first to analyze the effects of constant (over an orbit) body-fixed forces in the various yaw modes. Figure 3 summarizes these results in terms of the rate of change in semi-major axis when the nominal yaw mode switch points are utilized (Eq. 1). The analysis shows that semi-major axis is affected by forces in the X-direction only in the two fixed-yaw modes; orbital boost occurs when flying forward, and orbital decay when flying backward. During yaw steering the net forces in the X direction are essentially zero. Similarly, Y-direction forces are unimportant during the fixed yaw modes, but become important in the yaw steering modes. In positive yaw steering a +Y-axis force causes orbital decay at a rate that increases with β' angle. In negative yaw steering orbital boost of the same magnitude is expected. These distinctive signatures of X- and Y-body forces are helpful in interpreting the observed changes in semi-major axis.

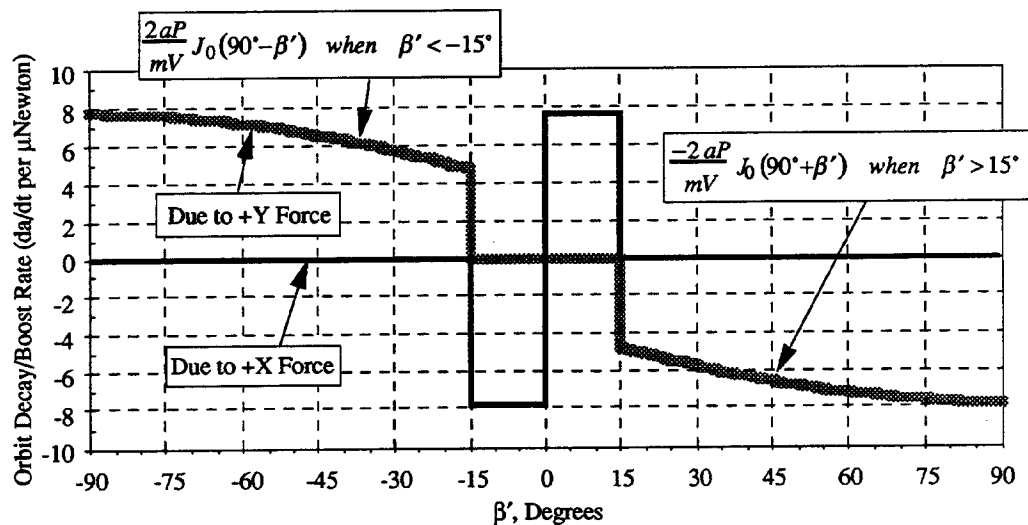


Figure 3 Daily Changes in da/dt (cm/day) vs β' for a Continuous Body-Fixed Thrust of $1\mu\text{N}$
(assumes nominal fixed yaw intervals: $-15^\circ < \beta' < 15^\circ$)

UNMODELED FORCES ARISING FROM THE SOLAR ARRAY OFFSET

The SA offset angle B (Eq. 3) is expected to be adjusted annually to compensate for performance degradation. A perturbation analysis was used to isolate the effects of reflections from the offset SA, as these forces are not modeled in the navigation software. Earth albedo effects and lags in reaching thermal equilibrium were ignored.

The force normal to the solar array due to specular and diffuse reflection was resolved into satellite body-fixed coordinates as a function of the solar orbit angle (Ω) for the different yaw modes. The body-fixed coordinates were then mapped into an earth-centered rotating coordinate system. An orbital average of the forces along the satellite velocity vector and their Fourier coefficients were then determined.

During yaw steering the average acceleration along the velocity is zero, even during periods having solar occultation. However, during fixed-yaw modes the average acceleration b_0 is

$$b_0 = a_N \sin B \left[\frac{\sin(\Omega_{occ}/2)}{\pi} \right] \quad (5)$$

for the *fly-forward* mode, where a_N is the magnitude of the acceleration due to the reflections from the solar array and Ω_{occ} is the angular length of the occultation interval. In the *fly-backward* mode these accelerations are in the opposite direction. If the SA offset is changed from lead to lag the force also reverses. All fixed-yaw periods to date have used a solar array *lead* angle (except the first two when $\psi_s = 0^\circ$). Otherwise, this X-direction force reverses at the yaw-flip. The acceleration a_N is

$$a_N = \frac{2G_S A_{sa}}{mc} \eta^* \cos^2 B \cos^2 \beta' \quad (6)$$

where G_S is the direct solar flux in W/m^2 , A_{sa} is the solar array area, η^* is a composite array reflectivity, and c is the speed of light. Combining Eqs. 5 and 6 and expressing Ω_{occ} in terms of β' ,

$$b_0 = \frac{2G_S A_{sa}}{mc} \eta^* (\sin B \cos^2 B) \left(\frac{\cos^2 \beta'}{\pi} \sqrt{1 - (\cos \lambda / \cos \beta')^2} \right), \quad |\beta'| < \lambda \quad (7)$$

where $\lambda = 55.7^\circ$, the value of β' when occultations cease. The maximum value of b_0 occurs when $B = 35.3^\circ$. The decrease in b_0 is $\sim 8\%$ at $\beta' = 15^\circ$, and $\sim 31\%$ at 30° . Typically, b_0 is $\sim 1.26 \text{ nm/s}^2$ when $\beta' = 0^\circ$ and $B = 57.5^\circ$, predicting a boost rate in the mean semi-major axis of $\sim 23 \text{ cm/day}$. Eq. (7) describes the fly-forward mode; the acceleration reverses direction for the fly-backward mode, resulting in an orbital decay of similar magnitude.

Corroboration of our analyses has been provided by Richter's thermal analysis¹⁷ which reproduces observed solar array temperature quite well over the full range of β' angles.

ESTIMATING THE ANOMALOUS FORCE

By 21 September 1992 the maneuver sequence designed to acquire operational orbit conditions had been successfully completed.⁴ Each of the six maneuvers repeatedly interrupted newly-stable orbit conditions, precluding opportunities to confidently establish a trend in the anomalous along-track force. Fortunately, this limitation had been acceptable, since orbit changes induced by these maneuvers were much greater than the effects of the observed decay

phenomena. However, once operational orbit conditions had been achieved, effective ground track maintenance then depended on reliable estimates of the anomalous forces, since the related orbital decay and boost rates were generally the same order-of-magnitude as the expected decay due to atmospheric drag.

Plans were made to estimate an along-track thrust multiplier ($1 + \tau$) as part of routine orbit determination. Only τ would be estimated, and not the drag multiplier ($1 + \rho_1$) since these along-track forces could not be satisfactorily estimated simultaneously. Brief studies quickly established that a single τ acting over a tracking arc of at least five days was necessary for confident estimation. Shorter tracking arcs, or more frequent (daily) τ estimates, were generally poorly determined and unsuitable for trend analysis leading to a reliable prediction model.*

A strategy was adopted to estimate a single τ acting over a seven-day tracking arc. This technique establishes an average along-track thrust force referred to the center of the tracking data arc. A daily moving average τ was obtained by advancing the seven-day tracking arc by a day and dropping off the first day, so adjacent solutions were always based on six days of common tracking data. This method produces a reasonably smooth and consistent daily history of τ which can then be expressed as an equivalent rate of change in orbit mean semi-major axis, da/dt ,¹⁸

$$T_{LT} = (1 + \tau), \text{ in } \mu\text{N}$$

$$\text{then } \frac{da}{dt} = \frac{P}{\pi} \left(\frac{T_{LT}}{m} \right) \equiv 7.7 T_{LT} \left(\frac{da}{dt} \text{ is in cm / day when } T_{LT} \text{ is in } \mu\text{N} \right) \quad (8)$$

where T_{LT} is the along-track thrust, τ is the estimated thrust parameter, P is the orbit period, and m is the satellite mass.

Finally, orbit determination solutions were isolated into distinct families corresponding to each of the four yaw control modes. This grouping technique prevents corruption of an orbit solution with dramatically different along-track force levels known to be present in adjacent yaw control modes (e.g., Figure 3). Ultimately, the desired prediction model should cover a period of at least three to four months, the typical spacing between orbit maintenance maneuvers. Such a model would define the behavior of the anomalous force during each of the four yaw control modes, as each mode would be used at least once between maneuvers.

The anomalous forces are first examined during periods of negative and positive yaw steering, then for fixed yaw, flying both forward and backward. Finally, these individual segments are assembled to form a composite prediction model for use orbit maintenance.

FORCES OBSERVED DURING NEGATIVE YAW STEERING

Stable orbit conditions suitable for estimating the anomalous force were first available in late September 1992, soon after achieving the operational orbit. The satellite was in a period of negative yaw steering ($\beta' < 0$). Unfortunately, this initial opportunity was necessarily brief, since the first orbit maintenance maneuver (OMM1) had already been planned to reverse the eastward drift of the ground track and to begin the orbit maintenance phase. Recognizing that predictions of the anomalous force were now needed for effective OMM1 design, the Project permitted the ground track to drift beyond the eastern control boundary, scheduling OMM1 on 12 October 1992 at the transition between Cycles 2 and 3 (originally OMM1 would have occurred ~10 days earlier between Cycles 1 and 2).**

* Early in the mission two 40-min passes/orbit of two-way Doppler via TDRS East and West were used; in May 1993 tracking was voluntarily reduced to one 40-min pass/orbit of one-way Doppler, a strategy consistent with the pre-launch metric tracking data acquisition plan.

** See the ground track history shown later in Figure 12.

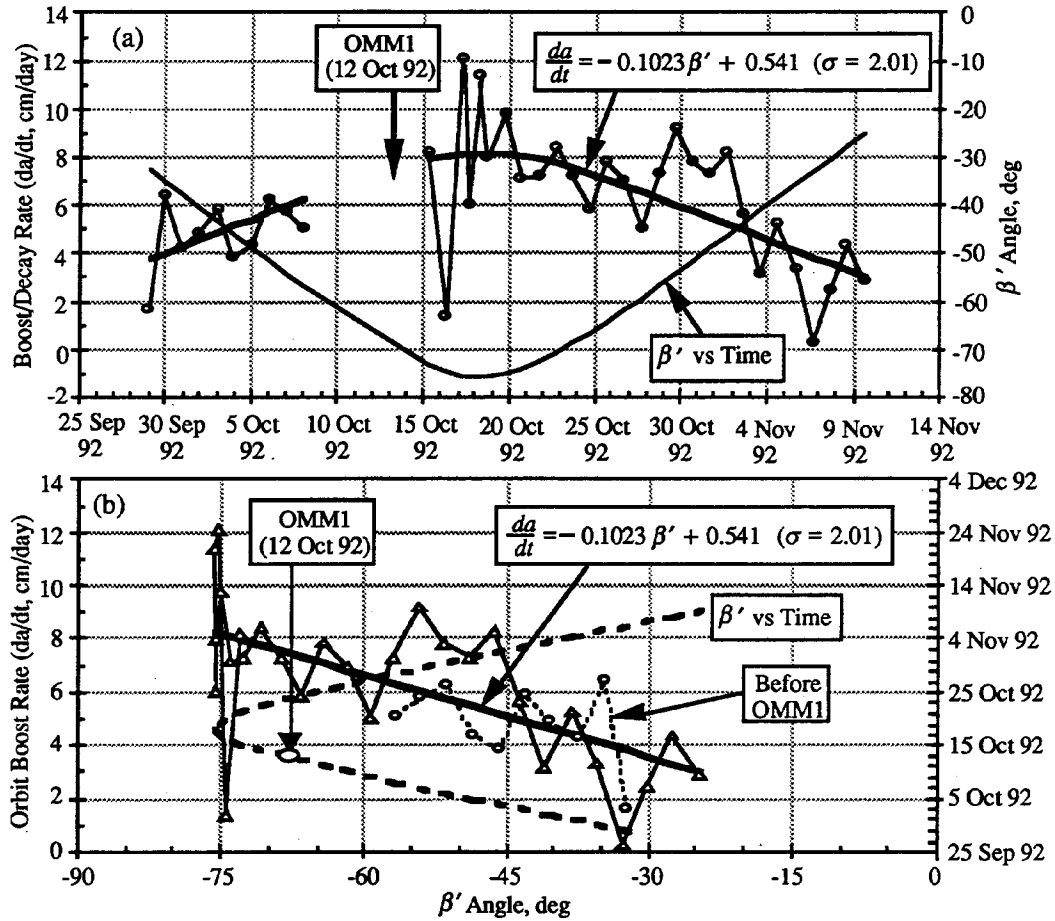


Figure 4(a) Estimate of Orbital Boost Due to Anomalous Force During Negative Yaw Steering (Time Dependence), and (b) β' Dependence

This operational plan provided a limited three-week period prior to OMM1 to obtain initial estimates of the anomalous force and to develop a post-maneuver prediction model. Figure 4(a) shows that a small orbital boost rate was observed prior to OMM1, but these data did little toward development of a meaningful post-maneuver prediction model. Quick review of Figure 2(a) during the period prior to OMM1 suggests that outgassing might have been the only force present, and may be nearing the end of an exponential decay. Because results from the daily τ estimates were so inconclusive, a decision was made to apply a continuation of the exponential trend for an additional 17 days after OMM1, believing this would account for the remaining outgassing and possibly mark the end of any anomalous forces. The modest consequences of this incorrect choice in relation to the overall effectiveness of subsequent ground track maintenance are discussed later.

Continued τ estimation after OMM1 (Figure 4a) shows sustained orbit boost rates between ~ 3 and ~ 8 cm/day that are negatively correlated with the β' variation. A least-squares fit of the daily da/dt values over the entire negative yaw steering period exhibits a linear variation with β' , as shown by Eq. (6) and illustrated in Figure 4(b).

$$\frac{da}{dt} = -0.1023\beta' + 0.541 \left(\frac{da}{dt} \text{ in cm / day when } \beta' \text{ in deg} \right) \quad (9)$$

FORCES OBSERVED DURING POSITIVE YAW STEERING

The anomalous force was first estimated during positive yaw steering in late Nov 92. This interval was briefly interrupted by OMM2 on 21 December 1992. Daily τ estimates indicated sustained orbital decay (Figure 5a) which were negatively correlated with the β' variation, a trend similar to that previously observed during negative yaw steering. Figure 5(b) and Eq. (10) show the orbital decay rate varies linearly with β' .

$$\frac{da}{dt} = -0.1023\beta' - 3.131 \quad (\beta' \text{ in deg, } da/dt \text{ in cm/day}) \quad (10)$$

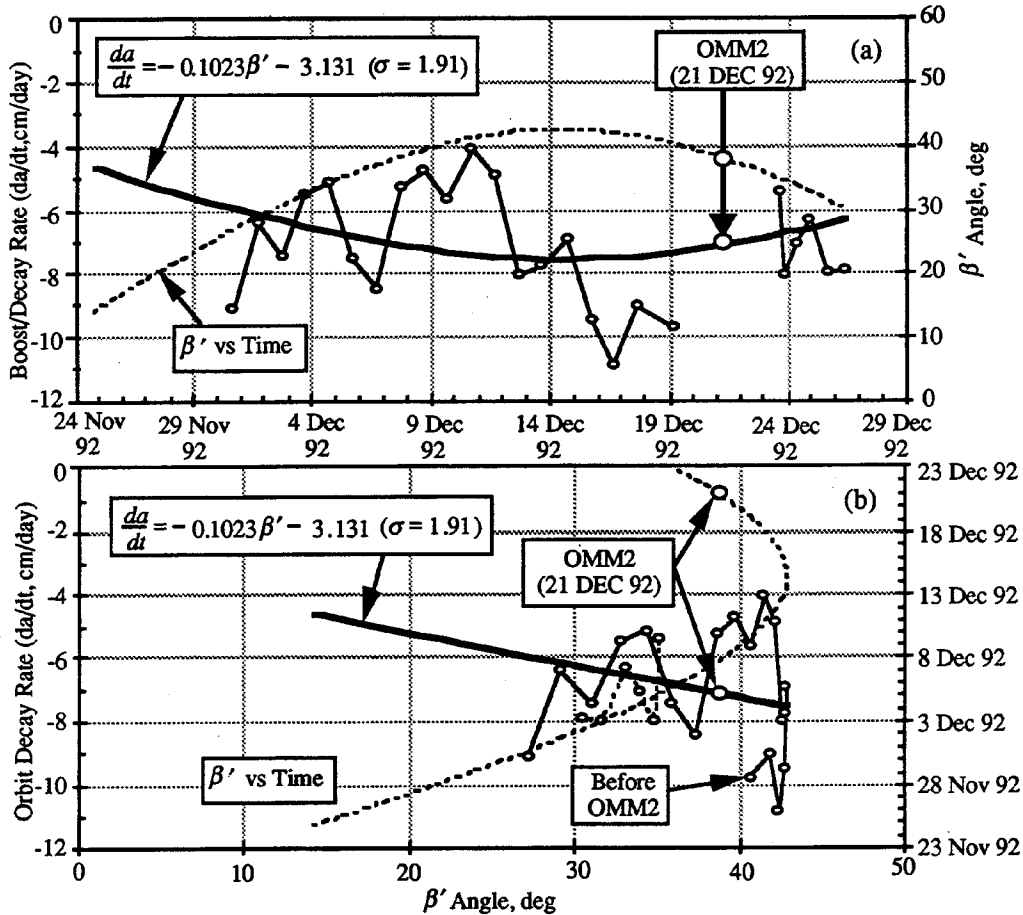


Figure 5(a) Estimate of Orbital Decay Due to Anomalous Force During Positive Yaw Steering (Time Dependence), and (b) β' Dependence

While the trends in the anomalous force with β' were similar, the observed orbital decay levels are ~ 3.7 cm/day greater than were predicted by variations with β' determined earlier during negative yaw steering. A comparison of Eqs. 9 and 10 clearly describes this difference (also see Appendix A). Marshall¹⁵ confirms with precision orbit determination using satellite laser ranging that the +Y forces were persistently different in these negative and positive β' regimes. Differences in orbit-sun geometry may explain these differences. The average along-track force due to SA curling should be smaller during occultations due to the reduced time in sunlight.

When $|\beta'| \geq 55.7^\circ$, the satellite orbit enters and remains in full sun as β' varies through its peak value.* The peak β' value in October 1992 was $\sim 76^\circ$; the orbit was in full sun for ~ 20 days. However, during the winter solstice in December 1992 the peak β' value reached only $\sim 42^\circ$ so full sun conditions were never achieved. This geometry repeats during every winter and summer solstice, while all the other β' cycles always provide extended periods of full sun. Richter's results¹⁷ also predict an increase in the +Y force upon entering full sun.

EMPIRICAL MODELS FOR POSITIVE AND NEGATIVE YAW STEERING

The previous sections have described the technique for estimating the anomalous force during positive and negative yaw steering modes, using two early sample periods. Figures 4 and 5, and Eqs. 9 and 10, summarize these results. Through June 93, there have been a total of five different periods of yaw steering, three in the negative and two in the positive yaw steering mode.

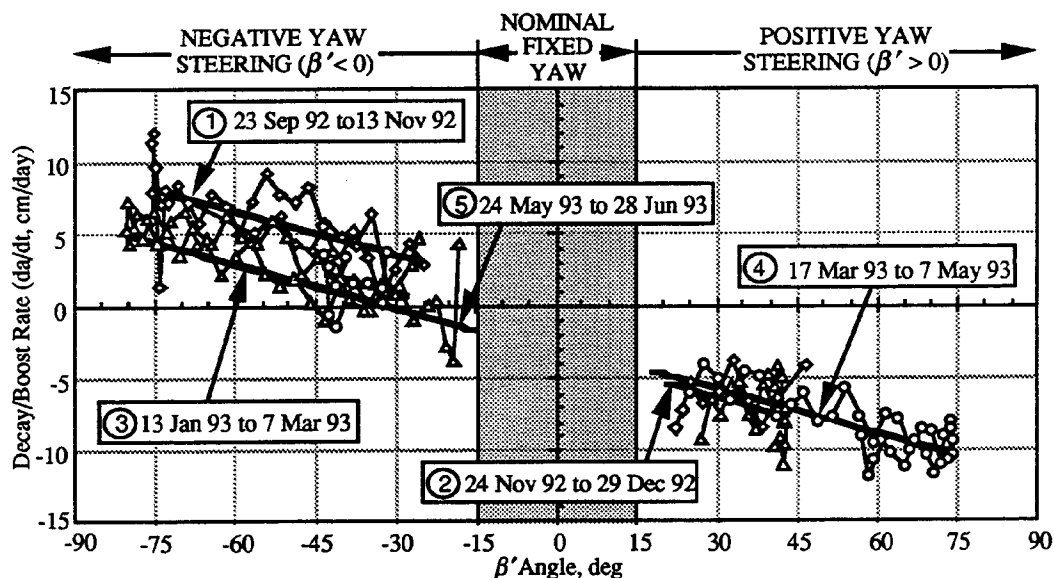


Figure 6 Composite Estimate of Orbital Decay/Boost Due to Anomalous Force for Both Positive and Negative Yaw Steering (β' Dependence)

Figure 6 shows the respective orbital boost or decay rates with β' for each of the five yaw steering periods. Periods 1, 3, and 5 are consecutive repeats of the negative yaw steering mode when an orbital boost was observed. The boost rates estimated for period 1 are ~ 3.7 cm/day higher than those for period 3, whereas period 3 is ~ 0.3 cm/day higher than period 5. The large reduction in boost rates between periods 1 and 3 may indicate additional decay due to residual outgassing, as the observed change is much larger than could be explained by errors in estimating τ . Also, periods 1 and 3 are intervals of full sun, making direct comparisons more valid. Period 5 occurs near the 1993 summer solstice, and so it is reasonable to expect it to differ somewhat from the trend observed during period 3.

The two intervals of positive yaw steering shown in Figure 6, periods 2 and 4, result in orbital decay. The empirical models for these two intervals indicate that period 4 has ~ 0.5 cm/day less decay than period 2, a trend which supports the presence of additional outgassing. However, period 4 includes an interval of full sun, while period 2 does not, so direct comparison of these two periods may be less valid.

* See the β' history since launch presented later in Fig. 11.

Overall, the anomalous forces during the four most recent yaw steering periods have exhibited excellent repeatability, certainly well within the estimation accuracy of the empirical models (see Appendix A). The remaining task prior to establishing an overall prediction model suitable for use in orbit maintenance maneuver design requires estimates of the anomalous force during periods of fixed yaw.

FORCES OBSERVED DURING FIXED YAW

Periods of fixed yaw nominally occur when $-15^\circ < \beta' < 15^\circ$. Five have occurred since launch. As indicated in Eq. (1) there are two fixed yaw control modes: flying forward when $0^\circ < \beta' < 15^\circ$, and flying backward when $-15^\circ < \beta' < 0^\circ$. These modes are always separated by a yaw flip maneuver near $\beta' = 0^\circ$.

An empirical model of the anomalous force was also developed for the two fixed-yaw modes by estimating a single thrust parameter τ over the tracking data arc.* The tracking data were carefully limited to the period of fixed yaw to avoid corruption of τ by data from adjacent yaw control modes where the along-track forces are dramatically different. A single fixed yaw mode covers about five days when nominal switch points are utilized. Since confident recovery of τ requires at least five days of tracking data, these solutions were sometimes marginally adequate. While the average force level could be estimated, it was usually not possible to establish any trend which might be related to β' variations during fixed yaw.

However, there have been two periods of fixed yaw extending outside $-15^\circ < \beta' < 15^\circ$ to support special circumstances. These longer periods of fixed yaw provided additional tracking data to facilitate more confident recovery of τ , while also presenting unique opportunities to observe any trend in the anomalous force during fixed yaw at β' angles above the nominal 15-deg limit. The first such occurrence was between 29 December 1992 and 8 January 1993 (~10 days). During this holiday period the satellite was placed in fixed yaw early when $\beta' \approx +23.7^\circ$ while flying forward, remaining in this mode until a yaw flip was performed near $\beta' = 0^\circ$. The second occurrence was between 13 and 24 May 93 (~11.5 days) when the satellite remained in fixed yaw (flying backward) longer, until $\beta' \approx -26.7^\circ$, in order to induce additional orbital decay to prevent the satellite ground track from drifting outside the western control boundary soon after executing OMM3 on 30 March 1993.**

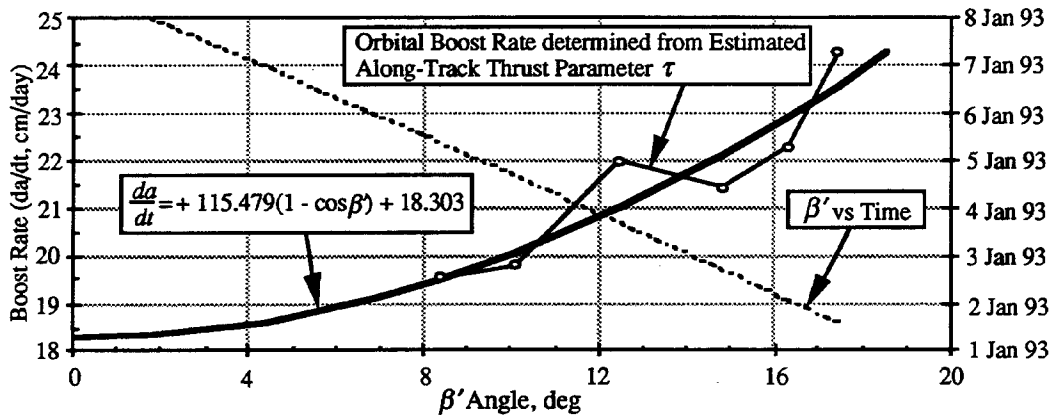


Figure 7 Orbital Boost Rate During Fixed Yaw and Positive β' ($\psi_s = 0^\circ$) from 29 Dec 92 to 8 Jan 93 (~10 days)

* Tracking data acquisition during fixed yaw is the same as previously described for yaw steering (see p. 9).

** See the ground track history shown later in Figure 12.

Figure 7 shows an orbital boost that increases with β' during the period of fixed yaw between 29 December 1992 and 8 January 1993. The empirical model (Appendix A) describes this variation as a simple linear function of $(1 - \cos\beta')$. However, this result does not support intuition that the along-track forces would decrease as β' increases. Such a result would be expected as the solar incidence angle increases with β' , thereby reducing the SA area projected normal to the along-track direction. In contrast, this surprising result was not duplicated in May 1993 while flying backwards in fixed yaw. Figure 8 shows the +X forces reverse to cause orbital decay, steadily declining as β' increases, while varying linearly with $(1 - \cos\beta')$, as also shown by the empirical model in Appendix A.

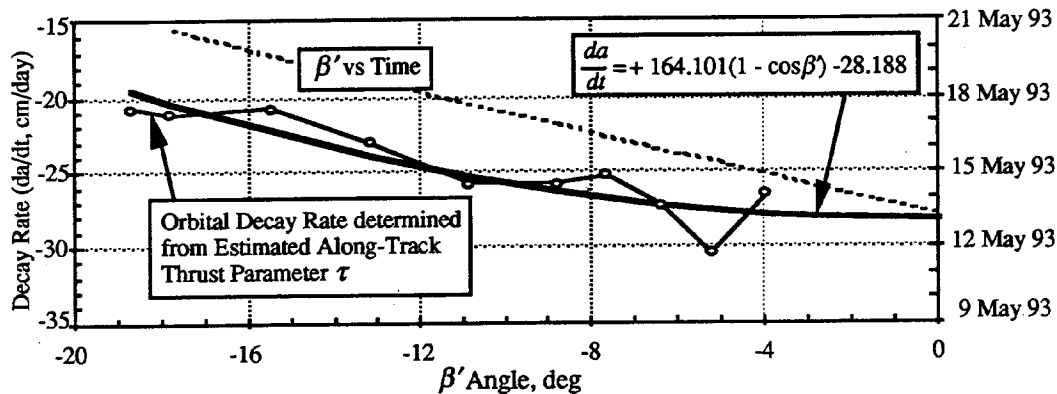


Figure 8 Orbital Decay Rate During Fixed Yaw and Negative β' ($\psi_s = 180^\circ$) from 13 to 24 May 1993 (~11.5 days)

Before confronting the obvious conflict presented by these contradictory results, let us first compare them with da/dt estimates obtained during other periods of fixed yaw. Figure 9 combines the results of Figures 7 and 8 with estimates from other fixed yaw periods and these comparisons are quite good, but are necessarily limited to β' values near 7.5° , the average β' during nominal periods of fixed yaw covering about five days.

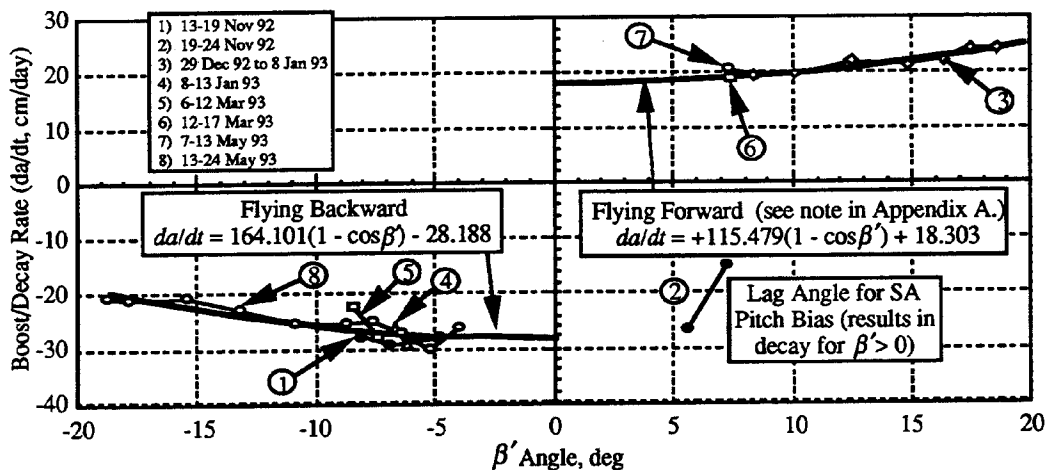


Figure 9 Orbital Boost/Decay Rates During Fixed Yaw for Positive and Negative β'

Richter's analysis¹⁷ supports our analytical results (Eq. 7) that the orbital boost declines as β' increases. Also, quick-look orbit determination performed by the JPL TOPEX/Poseidon Verification Team using laser ranging measurements confirm this trend. Ironically, the empirical function derived directly from τ better describes the expected variation if the term $+115.479(1 - \cos\beta')$ shown in Figure 7 were replaced by its negative: $-115.479(1 - \cos\beta')$, as is shown in Appendix A. Use of this more plausible model is planned for predicting future variations in da/dt with β' while flying forward in fixed yaw beyond $\beta' = -15^\circ$. Additional experience in this control mode will help establish the appropriate empirical model.

COMPOSITE EMPIRICAL MODEL

Observations of the anomalous force during each yaw control mode have been used to develop empirical models of the rate of change in the mean semi-major axis (da/dt) as a function of the sun-orbit plane separation angle β' . These results were derived from daily estimates of an average along-track thrust parameter τ obtained from moving seven-day tracking arcs. Appendix A summarizes the resulting empirical models, identifying the time interval and β' values for each yaw mode.

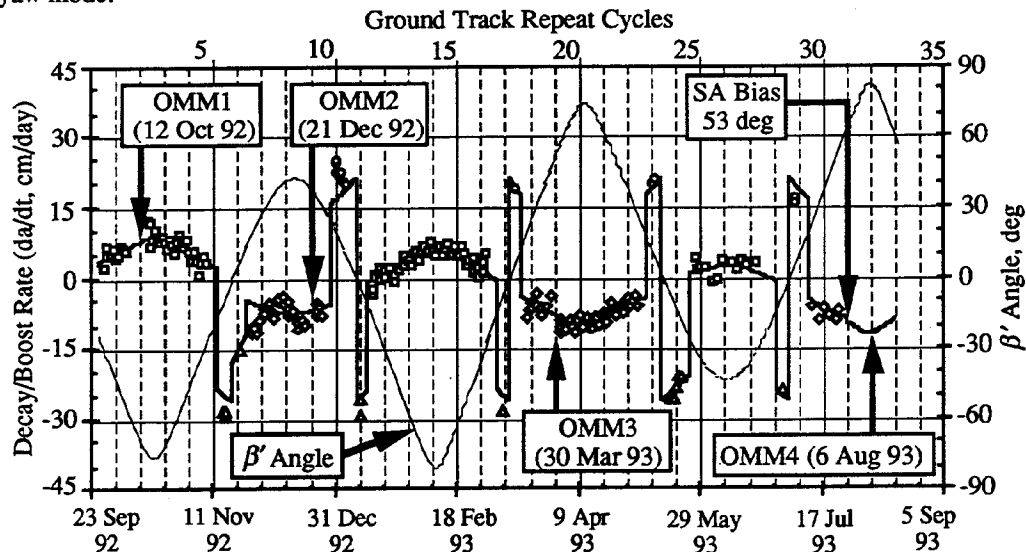


Figure 10 History of Orbital Decay and Boost Rates due to the Anomalous Force

Figure 10 shows a composite empirical model for late September 1992 through mid-August 1993.* This figure also displays the individual daily estimates of da/dt and locates each maintenance maneuver performed since launch. The variation in β' helps identify when yaw mode changes occur. During fixed yaw orbital boost occurs when $\beta' > 0^\circ$, orbital decay when $\beta' < 0^\circ$, while the *sawtooth* discontinuity results from yaw flips performed near $\beta' \approx 0^\circ$,

The results described by the empirical model during periods of yaw steering are consistent with a force acting along the body +Y axis (Figure 3). These forces appear to result from curling of the SA induced by a thermal gradient between the front and rear surfaces (Richter¹⁷). During fixed yaw the forces are generally consistent with the predicted effects of the SA pitch offset angle B summarized by Eq. (7).

* The discontinuity in the empirical model between orbit cycles 31 and 32 reflects the expected change induced by reducing the SA pitch offset angle B from 57.5 to 53 deg.

EFFECTS ON THE SATELLITE GROUND TRACK

Prior to launch, atmospheric drag was believed to be the most significant perturbation affecting the satellite ground track. Luni-solar gravity has predictable and significant periodic effects on the ground track comparable to those of drag,⁶ and is also routinely modeled as part of ground track maintenance. The relative effects of errors in predicting the anomalous forces and atmospheric drag are compared in Figure 11. Observed prediction errors are defined for the intervals between each orbit maintenance maneuver (Figure 11a). A comparison of these three curves shows progressively better modeling accuracy since the anomalous force was initially observed.

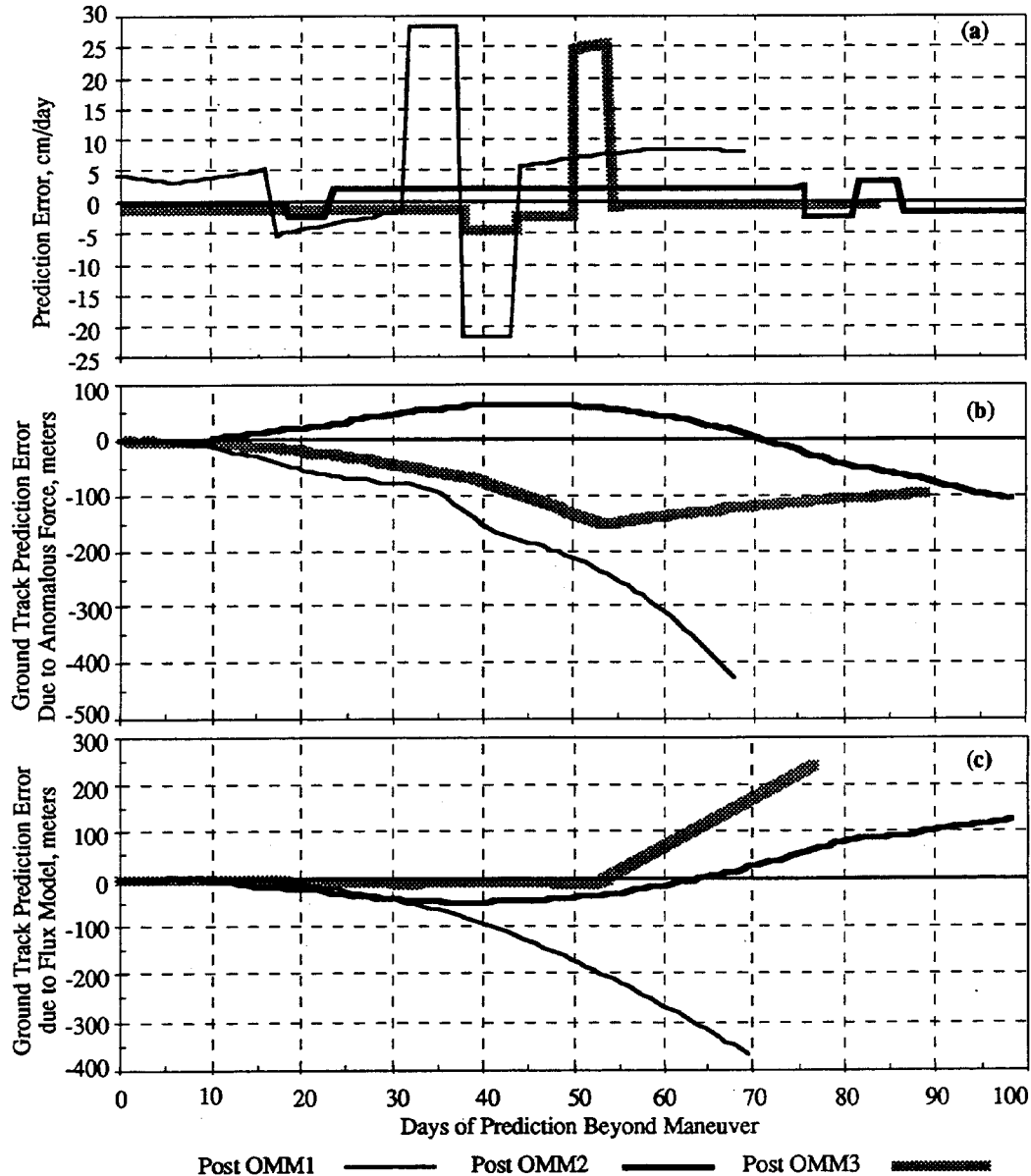


Figure 11(a) Error in Prediction of Anomalous Force, (b) Effect upon the Ground Track, and (c) Error in Ground Track Prediction due to Errors in the Flux Prediction Algorithm.

The translation of these prediction errors into ground track units is shown in Figure 11(b), and the corresponding ground track errors due to solar flux prediction errors are shown in Figure 11(c). These ground track errors were generated utilizing an analytic mean element propagation program (GTARG) which models all perturbations known to cause significant variations in the satellite ground track.¹⁹ These perturbations include earth oblateness (zonals to J_{20}), luni-solar gravity, atmospheric drag, and the anomalous force. Drag predictions were obtained by repeating the NOAA/SESC 27-day outlook²⁰ as required to cover the prediction interval.⁷ The ground track plots of Figure 11 were generated by differencing GTARG runs, but varying the boost (Figure 11b) and drag forces (Figure 11c).

Physically, we expect that an increase in the semi-major axis (from boost) will cause the ground track to accelerate westward, while a decrease moves the ground track eastward. Examination of Figure 11 confirms this physical intuition. If either the anomalous force or drag causes a smaller than expected decay, or the anomalous force causes a larger than expected boost, the true ground track falls west of the predicted ground track. These variations lead to the positive errors shown in Figure 11. Similarly, if the decay due to drag is more than expected, or the boost from the anomalous force is less than expected, the true ground track falls east of the prediction and this results in negative errors. For example, models available at the time of OMM1 overestimated the initial predicted boost and ignored it altogether after 17 days. Initial overestimation led to an increasingly negative ground track error; mismodeling after 17 days and omitting the large effects of the fixed-yaw periods led to large errors in the expected directions.

Significant modeling improvements were realized by the time of the OMM2 design. Slightly less than the expected decay was observed initially and during the first fixed yaw period, leading to a positive ground track prediction error. For approximately the next 50 days, the net boost during yaw steering was overestimated, causing the ground track error to move continuously less westward and eventually turn eastward. These observations and the increased predictability of the anomalous force led to the proposal to utilize the anomalous force for ground track maintenance by adjusting the duration of the fixed yaw periods. This strategy was used during the first fixed yaw period following OMM3, as the ground track approached the western boundary, thereby obviating the need for an additional orbit maintenance maneuver.²¹ As shown by Figure 12, our improved modeling has led to an increasing duration between maneuvers.

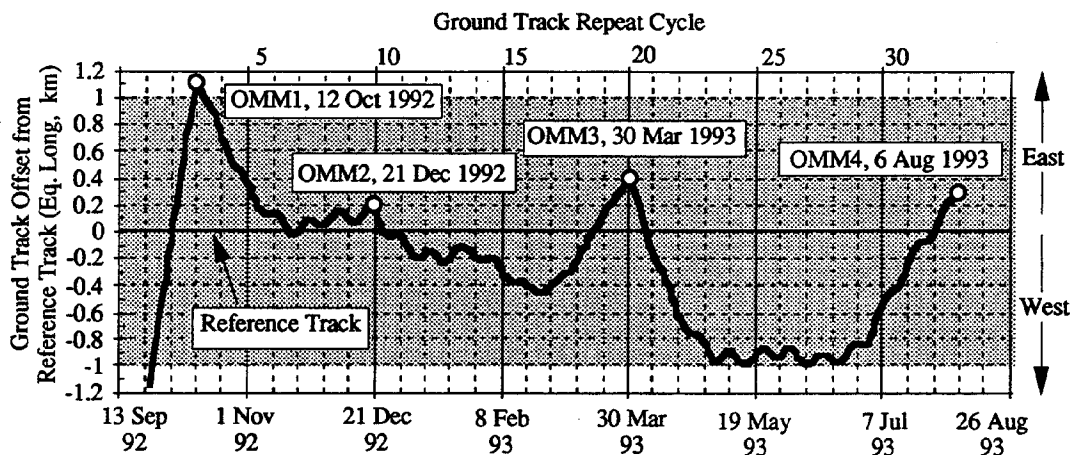


Figure 12 TOPEX/Poseidon Ground Track History (the Control Band is shaded).

SUMMARY

This paper describes the process by which the NAVT and the GSFC/FDF first detected, then began estimating, and finally developed an effective empirical model of the anomalous body-fixed forces. The primary purpose of this task was to develop a prediction model for use in ground track maintenance. This objective has been realized by achieving increasingly longer times between orbit maintenance maneuvers. The spacing between OMM1 and OMM2 was 70

days, 100 days between OMM2 and OMM3, and 130 days between OMM3 and OMM4 (Figure 12).

Reasonable physical interpretations of the forces have also been identified, although further refinements are expected and will be sought. Future work may include improved modeling of the changing SA curl effects with β' during yaw steering, particularly during full sun. Expected modeling changes imposed by adjusting the SA pitch bias are better understood and now must be verified and adjusted to reflect flight data.

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Last, but certainly not least, we owe great thanks and appreciation to the many GSFC/FDF personnel, and Ed Doll in particular, for their very professional and timely operational orbit determination support.^{22,23}

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Appendix A.
Empirical Models for the Anomalous Force vs Yaw Control Mode

YAW CONTROL MODE	START (UTC)	END (UTC)	EMPIRICAL MODEL (da/dt, cm/day, deg)	σ cm/d
Yaw Steering, $\beta' < 0$	23 Sep 92, 22:24:54	13 Nov 92, 08:00:08	$-0.1023\beta' + 0.541$	2.01
Fixed Yaw($\psi_s = 180^\circ$) $-14.8^\circ < \beta' < 0.2^\circ$	13 Nov 92, 08:28:18	19 Nov 92, 01:45:57	-28.41	0.65
Fixed Yaw($\psi_s = 0^\circ$) $+0.2^\circ < \beta' < +14.1^\circ$	19 Nov 92, 02:42:11	24 Nov 92, 18:08:05	-14.93	1.5
Yaw Steering, $\beta' > 0$	24 Nov 92, 18:36:12	29 Dec 92, 18:38:03	$-0.1023\beta' - 3.131$	1.91
Fixed Yaw($\psi_s = 0^\circ$) $+23.6^\circ < \beta' < -0.3^\circ$	29 Dec 92, 19:06:10	8 Jan 93, 18:34:24	$-115.479(1 - \cos\beta') + 18.303$ (see footnote*)	0.70
Fixed Yaw($\psi_s = 180^\circ$) $-0.3^\circ < \beta' < -13.3^\circ$	8 Jan 93, 19:30:22	13 Jan 93, 17:07:42	-27.53	2.26
Yaw Steering, $\beta' < 0$	13 Jan 93, 17:35:50	7 Mar 93, 00:12:30	$-0.1023\beta' - 2.228$	1.56
Fixed Yaw($\psi_s = 180^\circ$) $-15^\circ < \beta' < 0^\circ$	7 Mar 93, 00:40:37	12 Mar 93, 02:57:59	-25.22	3.86
Fixed Yaw($\psi_s = 0^\circ$) $0^\circ < \beta' < 15^\circ$	12 Mar 93, 03:54:00	17 Mar 93, 08:04:05	+18.6	1.48
Yaw Steering, $\beta' > 0$	17 Mar 93, 08:32:18	7 May 93, 12:51:02	$-0.1023\beta' - 2.560$	1.48
Fixed Yaw($\psi_s = 0^\circ$) $+15^\circ < \beta' < 0^\circ$	7 May 93, 13:19:04	13 May 93, 07:33:12	+20.04	0.84
Fixed Yaw($\psi_s = 180^\circ$) $0^\circ < \beta' < -26.5^\circ$	13 May 93, 08:29:10	24 May 93, 19:34:39	$164.101(1 - \cos\beta') - 28.188$	1.41
Yaw Steering, $\beta' < 0$	24 May 93, 20:03:04	28 Jun 93, 18:12:00	$-0.1023\beta' - 1.935$	1.84
Fixed Yaw($\psi_s = 180^\circ$) $-13.1^\circ < \beta' < 0^\circ$	28 Jun 93, 18:40:00	4 Jul 93, 02:36:00	-23.97	1.48
Fixed Yaw($\psi_s = 0^\circ$) $0^\circ < \beta' < 15^\circ$	4 Jul 93, 03:32:00	12 Jul 93, 03:11:00	+15.91	1.48

* A least-squares fit of da/dt determined from equivalent estimates of \mathcal{T} resulted in the term $+115.479(1 - \cos\beta')$ (see Figure 7), which implies an increase in da/dt with β' when a decrease would be expected. Richter's analyses¹⁷ and quick-look orbit determination by the JPL Verification Team using satellite laser ranging data confirm a decreasing trend. The term $-115.479(1 - \cos\beta')$ listed above better describes the expected variation with β' and will be used for predictions until additional flight data are available to refine the model.